

# Photomagnetism in Cyano-Bridged Bimetal Assemblies

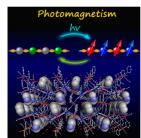
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## CONSPECTUS

The study of photoinduced phase-transition materials has implications for the fields of inorganic chemistry, solid-state chemistry, and materials science. Cyano-bridged bimetal assemblies are promising photomagnetic materials. Because cyano-bridged bimetal assemblies possess various absorption bands in the visible light region, their electronic and spin states can be controlled by visible light irradiation. Moreover, the selection of magnetic metal ions and organic ligands provide a way of controlling spin—spin interactions through a cyano bridge.



In this Account, we describe cyano-bridged bimetal assemblies developed in our laboratory.  $Cu^{II}_{2}[Mo^{IV}(CN)_{8}] \cdot 8H_{2}O$  (CuMo),  $Rb^{I}Mn^{II}[Fe^{III}(CN)_{6}]$  (RbMnFe), and  $Co^{II}_{3}[W^{V}(CN)_{8}]_{2} \cdot (pyrimidine)_{4} \cdot M_{2}O$ 

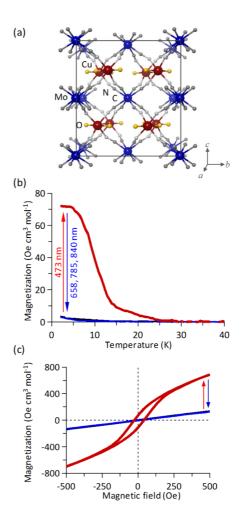
 $6H_2O$  (CoW) induce photomagnetism via photoinduced metal-to-metal charge transfers (MM'CT), while  $Fe^{II}_2[Nb^{IV}(CN)_8]$ . (4-pyridinealdoxime)<sub>8</sub> · 2H<sub>2</sub>O (FeNb) exhibits a photoinduced magnetization via a photoinduced spin crossover. Irradiation with 473 nm light causes the CuMo system to exhibit a spontaneous magnetization with a Curie temperature (*T<sub>c</sub>*) of 25 K, but irradiation with 532, 785, and 840 nm light reduces the magnetization. In this reversible photomagnetic process, excitation of the MM'CT from Mo<sup>IV</sup> to Cu<sup>II</sup> produces a ferromagnetic mixed-valence isomer of Cu<sup>I</sup>Cu<sup>II</sup>[Mo<sup>V</sup>(CN)<sub>8</sub>] · 8H<sub>2</sub>O (CuMo'). CuMo' returns to CuMo upon irradiation in the reverse-M'MCT band. RbMnFe shows a charge transfer (CT)-induced phase transition from the Mn<sup>III</sup> – Fe<sup>III</sup> phase to the Mn<sup>III</sup> – Fe<sup>III</sup> phase. Irradiation with 532 nm light converts the Mn<sup>III</sup> – Fe<sup>III</sup> phase into the Mn<sup>III</sup> – Fe<sup>III</sup> phase, and we observe photodemagnetization. In contrast, irradiation of the Mn<sup>III</sup> – Fe<sup>III</sup> phase with 410 nm light causes the reverse phase transition. A CT-induced Jahn – Teller distortion is responsible for this visible light-induced reversible photomagnetic effect. In the CoW system, a CT-induced spin transition causes the thermal phase transition from the Co<sup>III</sup> – W<sup>IV</sup> phase with 840 nm light causes ferromagnetism with a *T<sub>c</sub>* of 40 K and magnetic coercive field (*H<sub>c</sub>*) of 12 000 Oe, but excitation of the back M'MCT (Co<sup>II</sup> → W<sup>V</sup>) with 532 nm light leads to the reverse phase transition.

These examples of the photomagnetic effect have occurred by exciting MM/CT bands. In the fields of inorganic chemistry and materials science, researchers have studied extensively the photoinduced phase transitions between low-spin (LS) and high-spin (HS) transition metal ions. Recently, we have observed the first example of photoinduced spin crossover ferromagnetism with a FeNb system ( $T_c = 20$  K and  $H_c = 240$  Oe), in which a strong superexchange interaction between photoproduced Fe<sup>II</sup>(HS) and neighboring paramagnetic Nb<sup>IV</sup> operates through a CN bridge. The optical switching magnets described in this Account may lead to novel optical recording technologies such as optomagnetic memories and optical computers.

#### 1. Introduction

Research on photoinduced phase-transition materials is an attractive subject in the fields of inorganic chemistry, solid-state chemistry, and materials science.<sup>1–4</sup> Some of the applicative results have been utilized in industrial applications, including optical recording materials such as digital versatile discs (DVDs) and Blu-ray discs.<sup>5,6</sup> In addition, magnetic materials have a longer industrial history as magnetic recording materials in tapes and discs. For example, magneto-optical disks, in which recording is performed by heating using light a magnetic material beyond its Curie temperature ( $T_C$ ), are still being utilized to this day. Unlike these conventional magnetic recording materials, photomagnetic materials (optical switching magnets) described herein switch their physical properties when exposed to light, thereby leading to the potential development of novel optical recording technologies such as opto-magnetic memories and optical computers. Photomagnetic materials, therefore, have been receiving not only scholarly attention from the perspective of nonequilibrium phase transition, but also extensive industrial attention.

As candidate materials that demonstrate photomagnetism, metal assembled complexes are thought to be promising. Since metal complexes possess various absorption bands in the visible light region, there is a possibility of controlling their electronic states and spin states using visible light. Moreover, magnetic metal ions and organic ligands can be selected in view of spin-spin interactions, and thus material design suitable for each research objective can be conducted. Additionally, when a cyanometalate is bonded to a transition metal, the ligand field of the cyanometalate, that is, the energy levels of the d-orbitals, is influenced. For example, in an octacyanometalate building block, it takes various coordination geometries, for example, square antiprism, bicapped trigonal prism, dodecahedron, resulting in that various crystal structures can be formed. Motivated by these characteristics, various cyano-bridged metal-complex magnetic materials with zero-dimensional high-spin clusters, one-, two-, and three-dimensional (3D) magnetic structures have been designed and synthesized.<sup>7–35</sup> By focusing on the various coordination structures of metal ions and the structural flexibility of cyano groups, several studies have reported on functional metal complex magnetic materials that respond to physical stimuli such as light,<sup>36–57</sup> pressure,<sup>58,59</sup> and electric field,<sup>60,61</sup> as well as chemical stimuli such as humidity<sup>62</sup> and gaseous adsorption.<sup>63,64</sup> In particular, we have developed several magnetic functional cyano-bridged metal complexes, for example, photomagnetic materials, <sup>38,39,46–56</sup> magnetic materials exhibiting two compensation temperatures,<sup>65</sup> ferroelectric ferromagnetic metal complex,<sup>60</sup> piezoelectric ferromagnetic material exhibiting magnetization-induced second harmonic generation (MSHG),<sup>66,67</sup> humidity and solvent vapor sensitive magnetic materials,<sup>62,64</sup> and so on. Herein, we describe cyano-bridged magnetic metal complexes that we have developed: (i)  $Cu^{II}_{2}[Mo^{IV}(CN)_{8}] \cdot 8H_{2}O$ , (ii)  $Rb^{I}Mn^{II}[Fe^{III}(CN)_{6}]$ , and (iii)  $Co^{II}_{3}[W^{V}(CN)_{8}]_{2} \cdot (pyrimidine)_{4} \cdot$ 6H<sub>2</sub>O, which induce photomagnetic phenomena via photoinduced charge-transfer transition, and (iv) Fe<sup>II</sup><sub>2</sub>[Nb<sup>IV</sup>(CN)<sub>8</sub>]. (4-pyridinealdoxime)<sub>8</sub>·2H<sub>2</sub>O, which exhibits photoinduced spin-crossover ferromagnet.

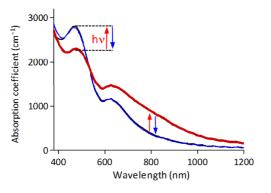


**FIGURE 1.** Visible-light-induced reversible photomagnetism in the CuMo octacyano-complex. (a) Schematic crystal structure. (b) Magnetization vs temperature curves before (black line) and after light irradiation with 473 nm light (red line), after irradiation with 658, 785, or 840 nm light (blue line). (c) Magnetic hysteresis loops at 3 K before (black line) and after irradiating with 473 nm light (red line), after thermal treatment of 250 K (blue line).

### 2. Visible-Light-Induced Reversible Photomagnetism of a CuMo Octacyano-Complex Based on Photoinduced Charge Transfer

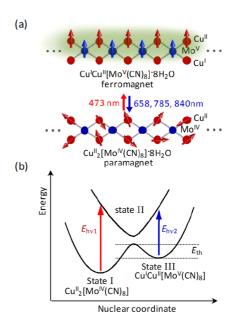
In mixed-valence metal complexes, metal-to-metal charge transfer (MM'CT) absorption bands are observed in the visible region. If charge transfer can be induced by the photoexcitation of the MM'CT, the electronic states of mixed-valence metal complexes can be changed. Herein, the visible-light-induced reversible magnetism of CuMo octacyano mixed valence metal complex,<sup>53,55</sup>  $Cu_2[Mo(CN)_8] \cdot 8H_2O$ , is described.

This complex, obtained as a purple powder, has a 3D structure where Cu and Mo are alternately bridged via a cyano group (Figure 1a). Magnetization vs temperature curves of this complex before and after irradiation are



**FIGURE 2.** UV–vis absorption spectra of the CuMo octacyano-complex at 3 K before (black line) and after irradiation with 473 nm light (red line) and after thermal treatment of 250 K (blue line).

shown in Figure 1b. Before irradiation, the paramagnetic behavior with  $Cu^{II}$  (S = 1/2) was observed. Upon exposure to a 473 nm continuous wave (cw) laser light, spontaneous magnetization was induced, resulting in a ferromagnetic state with  $T_{\rm C} = 25$  K. Magnetization vs magnetic field curve after irradiation indicated magnetic coercive field  $(H_c)$  of photoinduced phase is 30 Oe (Figure 1c). Conversely, when this photoinduced ferromagnetic phase was irradiated with laser light with  $\lambda \geq 520$  nm, we observed decreases in magnetization. Figure 1b depicts such decreases in magnetization upon light irradiation with a wavelength of 658, 785, or 840 nm. The UV-vis absorption spectrum of filmtype of Cu<sub>2</sub>[Mo(CN)<sub>8</sub>]·8H<sub>2</sub>O showed the MM'CT band from Mo<sup>IV</sup> to Cu<sup>II</sup> around 480 nm (Figure 2). After irradiation with 473 nm light, reverse-M'MCT band from Cu<sup>I</sup> to Mo<sup>V</sup> was observed around 710 nm. The photoinduced spontaneous magnetization is thought to be induced by the photoinduced charge-transfer transition from Mo<sup>IV</sup> (S = 0) to Cu<sup>II</sup> (S = 1/2). After the charge transfer was completed, the composition of its valence isomer became Cu<sup>I</sup>Cu<sup>II</sup>[Mo<sup>V</sup>(CN)<sub>8</sub>]·8H<sub>2</sub>O, with unreacted Cu<sup>II</sup> remaining. It is thought that a superexchange interaction was induced between the unreacted Cu<sup>II</sup> and the photoproduced  $Mo^V$  (S = 1/2) via cyano groups, which lead the ferromagnetism (Figure 3a). In contrast, photoinduced demagnetization indicates that the laser irradiations of 658, 785, or 840 nm caused the photoinduced back charge transfer from Cu<sup>I</sup> to Mo<sup>V</sup>, thereby restoring the initial state. It is worth noting that this photoinduced magnetization also disappeared upon a thermal treatment (250 K) and reverted to the initial state. The consideration of this reversible photomagnetic phenomenon in terms of the energy diagram of mixed-valence metal complexes<sup>68-71</sup> suggests that irradiation to MM'CT absorption bands with a 473 mn laser light produced a valence isomer. Irradiation to the back photoinduced charge transfer absorption bands with a



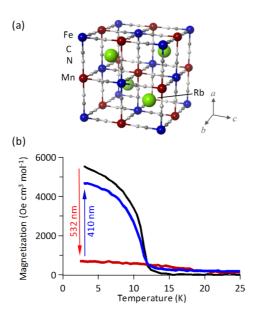
**FIGURE 3.** (a) Mechanism of magnetic ordering of photomagnetic phenomena in the CuMo octacyano-complex. (b) Mechanism of photoinduced charge transfer in a class II mixed-valence complex.

658, 785, or 840 nm laser light converted the valence isomer to the initial state with light (Figure 3b). This photomagnetic effect was repeatedly observed, indicating that this complex is a visible-light-induced, reversible photomagnetic material. Moreover, an analogue of this complex,  $Cs_2^I Cu_7^{II} [Mo^{IV}(CN)_8]_4 \cdot 6H_2O$ , can be electrochemically synthesized as a thin film, with similar photomagnetic properties.<sup>54</sup>

### 3. Visible-Light-Induced Reversible Photomagnetism of a RbMnFe Hexacyano-Complex Based on Photoinduced Charge-Transfer-Induced Jahn–Teller Distortion

Photomagnetism can be achieved by light irradiation onto magnetic materials which have a bistability. In bistable materials, the energy barrier between bistable states maintains a photoinduced state, resulting the observation of the persistent photoinduced magnetic state. From this perspective, RbMnFe hexacyano complex (Figure 4a),<sup>46–48,72–77</sup> one of a Prussian blue analogue, is suitable for observing photomagnetism. It is known to exhibit a charge-transfer phase transition from Mn<sup>II</sup> (S = 5/2)–NC–Fe<sup>III</sup> (S = 1/2) [high-temperature (HT)] phase to Mn<sup>III</sup> (S = 2)–NC–Fe<sup>III</sup> (S = 0) [low-temperature (LT)] phase accompanying Jahn-teller distortion on Mn<sup>III</sup>.

To excite an MM'CT band, the LT phase was exposed to a 532 nm laser light. Before light irradiation, magnetization vs temperature curve of the LT phase revealed that this phase



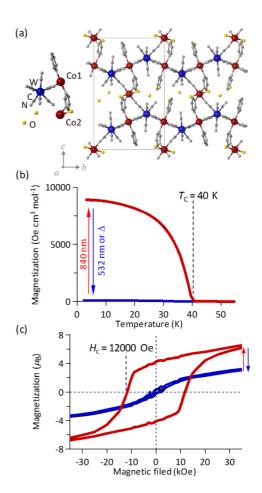
**FIGURE 4.** Visible-light-induced reversible photomagnetism in the RbMnFe hexacyano-complex. (a) Crystal structure. (b) Magnetization vs temperature curves before and after light irradiation. Before irradiation (black line), after irradiation with 532 nm light (red line), after irradiation with 410 nm light (blue line).

was ferromagnet with  $T_{\rm C} = 12$  K (Figure 4b). Additionally, upon exposure to a 532 nm laser light, a decrease in magnetization was observed, and upon subsequent exposure to a 410 nm laser light, an increase in magnetization was observed. Such a change phenomenon in magnetization was repeatedly observed by alternately irradiating the LT phase with 532 and 410 nm lasers.

This reversible photomagnetism can be explained by the following. Exposure to a 532 nm laser light excited the MM'CT (Fe<sup>II</sup>  $\rightarrow$  Mn<sup>III</sup>) band, inducing a photoinduced phase with the same electronic state as the HT phase. The relaxation of the metastable photoinduced phase to the stable LT phase was suppressed by thermal energy. In contrast, the ligand-to-metal charge transfer (LMCT) (CN<sup>-</sup>  $\rightarrow$  Fe<sup>III</sup>) band excitation of [Fe(CN)<sub>6</sub>]<sup>3-</sup> induced by irradiation with a 410 nm laser light resulted in the different photoexcited state, resulting in the transition to the initial LT phase. To our knowledge, this is the first observation of such an optical switching between a ferromagnetism and an antiferromagnetism.

### 4. Photoinduced Gigantic Coercivity of a CoW Octacyano-Complex Based on Photoinduced Charge-Transfer-Induced Spin Transition

Octacyano metal complexes can have various coordination structures that depend on the chemical environment, e.g., square antiprism  $(D_{4h})$ , dodecahedron  $(D_{2d})$ , and bicapped

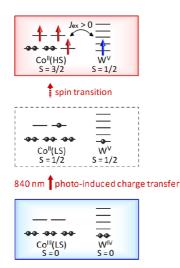


**FIGURE 5.** Crystal structure of the CoWprm octacyano-complex. (a) Asymmetric unit (left) and view along the *a*-axis (right). (b) Magnetization vs temperature curves before and after light irradiation. (c) Magnetic hysteresis loops before and after light irradiation (at 2 K). Before irradiation (black lines), after irradiation with 840 nm light (red lines), and after thermal treatment of 150 K or irradiation with 532 nm light (blue lines).

trigonal prism ( $C_{2v}$ ). Therefore, octacyano metal complexes have the potential to sensitively respond to external stimuli. Here we discuss photomagnetic phenomena exhibiting gigantic coercivity, which are observed in Co<sub>3</sub>[W(CN)<sub>8</sub>]<sub>2</sub>-(prm)<sub>4</sub>·6H<sub>2</sub>O (CoWprm complex),<sup>51,52</sup> a 3D structure containing pyrimidine (prm) as an organic ligand.

The crystal structure of a CoWprm complex, obtained as a red powder, belongs to the monoclinic system ( $P2_1/n$ ) where Co and W are alternatively bridged via cyano ligands to form a 3D network structure (Figure 5a). This complex showed a charge-transfer-induced spin transition as a function of temperature. The electronic state of the HT phase is  $Co^{II}(HS; S = 3/2)_3 - W^V(S = 1/2)_2$  and that of the LT phase is  $Co^{II}(LS; S = 0)_2 - W^{IV}(S = 0)_2 - Co^{II}(HS; S = 3/2)$ . This transition is induced by the MM'CT between Co and W and the spin transition in  $Co^{III}$ . Ultraviolet–visible (UV–vis) absorption spectra revealed a peak around 500 nm assigned to MM'CT,

 $Co^{II} \rightarrow W^V$ , in the HT phase and a peak around 750 nm assigned to M'MCT,  $W^{IV} \rightarrow Co^{III}$ , in the LT phase.

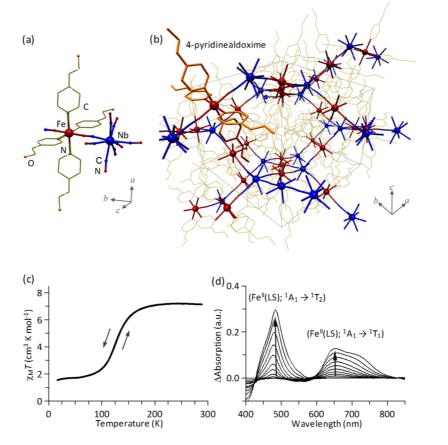


**FIGURE 6.** Mechanism of the photomagnetic phenomenon in the CoWprm octacyano-complex due to photoinduced charge-transferinduced spin transition.

The LT phase of the CoWprm complex is paramagnetic. Upon light irradiating ( $\lambda = 840$  nm), the sample exhibits spontaneous magnetization (Figure 5b).  $T_{\rm C}$  and  $H_{\rm c}$  of this photoinduced magnetic phase are 40 K and 12 000 Oe, respectively (Figure 5c). With the subsequent irradiation by a laser light ( $\lambda = 532$  nm), the magnetization was decreased. Such a change in magnetization was repeatedly observed by alternatively irradiating the CoWprm complex with 840 and 532 nm lasers. Investigation of the electronic states of the CoWprm complex using infrared (IR) absorption spectra revealed that the visible-light reversible photomagnetic phenomenon observed was induced by the reversible charge-transfer-induced spin transition between the LT phase,  ${\rm Co}^{\rm II}({\rm LS})_2 - {\rm W}^{\rm IV}_2 - {\rm Co}^{\rm II}({\rm HS})$ , and the photoinduced phase,  ${\rm Co}^{\rm II}({\rm HS})_3 - {\rm W}^{\rm V}_2$  (Figure 6).

### **5. Photoinduced Spin Crossover Ferromagnetism of a FeNb Octacyano-Complex**

Transition metal complexes (from  $d^4$  to  $d^7$ ) with octahedral coordination can transit from a high-spin (HS) state to a



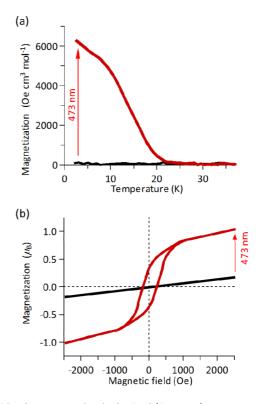
**FIGURE 7.** Crystal structure of the FeNb(4-pyoxm) octacyano-complex. (a) Asymmetric unit and (b) cyano-bridged Fe–Nb three-dimensional framework. (c)  $\chi_M T - T$  curves of the FeNb(4-pyoxm) octacyano-complex. (d) Temperature dependence of UV–vis differential spectra. Spectra are depicted from 300 to 100 K at 20 K intervals by setting a spectrum at 300 K as the standard.

low-spin (LS) as temperature decreases, which is called a "spin crossover transition". This phenomenon was first discovered in an Fe<sup>III</sup> complex, and many spin crossover complexes have been reported since then.78-80 Transition phenomena from the LS state to the HS state by light irradiation are termed "light-induced excited spin-state trapping (LIESST)" and have been studied as a technique to control paramagnetic states using light. Since its first discovery in an Fe<sup>II</sup> complex ( $[Fe^{II}(1-propyltetrazole)_6](BF_4)_2$ ),<sup>81</sup> a wide variety of studies have been conducted.<sup>78,80</sup> Spin crossover complexes that have been reported are mostly those composed of molecular crystals, where molecules gather to crystallize via the van der Waals force. However, if spin crossover complexes that have 3D network structures can be synthesized and HS sites can form a magnetic ordering with each other, then the spontaneous magnetization can be expected. From this perspective, we synthesized a new octacyano metal complex,  $Fe_2[Nb(CN)_8] \cdot (4-pyoxm)_8 \cdot 2H_2O$  (FeNb(4-pyoxm) complex), which contains Fe, Nb, and 4-pyridinealdoxime (4-pyoxm) as the organic ligand and then discovered ferromagnetism caused by photoinduced spin crossover,49 which is described below.

The crystal structure of an FeNb(4-pyoxm) complex is tetragonal (space group:  $I4_1/a$ ) as shown in Figure 7a and b, where an Fe<sup>II</sup> ion is connected to two N atoms of cyano groups and to four N atoms of 4-pyoxm. In addition, four cyano groups of Nb<sup>IV</sup>(CN)<sub>8</sub> are bridged with Fe ions, and the other four are not bridged. Fe and Nb ions are continually bridged via cyano groups to form a 3D network structure.

A measurement of the temperature dependence of the molar magnetic susceptibility ( $\chi_M$ ) of an FeNb(4-pyoxm) complex revealed a drastic decrease in the  $\chi_M T$  around 130 K (Figure 7c). In the UV–vis absorption spectra, absorption assigned to the d–d transitions ( ${}^{1}A_{1} \rightarrow {}^{1}T_{2}$ ,  ${}^{1}A_{1} \rightarrow {}^{1}T_{1}$ ) of Fe<sup>II</sup>(LS) (S = 0) that was absent at room temperature was observed around 480 and 650 nm at low temperature region (Figure 7d). Additionally, in the  ${}^{57}$ Fe Mössbauer spectra, Fe<sup>II</sup>(LS) was observed at low temperature. These results suggest that a decrease in  $\chi_M T$  around 130 K is induced by the spin crossover transition from Fe<sup>II</sup>(HS) (S = 2) to Fe<sup>II</sup>(LS) (S = 0).

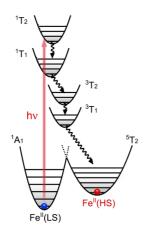
Next we discuss photomagnetic effects in an FeNb(4pyoxm) complex. Prior to light irradiation, the LT phase of an FeNb(4-pyoxm) complex is paramagnetic, while after light irradiation with a 473 nm cw laser light, spontaneous magnetization and magnetic hysteresis appear, suggesting that the paramagnetic phase changes ferromagnetic phase



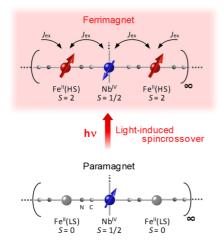
**FIGURE 8.** Photomagnetism in the FeNb(4-pyoxm) octacyano-complex due to photoinduced spin crossover. (a) Magnetization vs temperature curves before and after light irradiation. (b) Magnetic hysteresis loops before and after light irradiation at 2 K. Before irradiation (black lines) and after irradiation with 473 nm light (red lines).

(Figure 8). In the photoinduced ferromagnetic phase,  $T_{C}$ ,  $H_{C}$ , and saturation magnetization ( $M_{s}$ ) are 20 K, 240 Oe, and 7.4  $\mu_{B}$ , respectively. The  $M_{s}$  value of 7.4  $\mu_{B}$  agrees well with a calculated value of 7.7  $\mu_{B}$  obtained in a case where the spins of photoinduced Fe<sup>II</sup>(HS) (S = 2,  $g_{Fe^{II}(HS)} = 2.17$ ) and Nb<sup>IV</sup> (S = 1/2,  $g_{Nb^{IV}} = 1.99$ ) are antiparallel to each other. Before and after light irradiation with a 473 nm laser light, a UV–vis spectra at 3 K revealed a decrease in the absorption of  ${}^{1}A_{1} \rightarrow {}^{1}T_{2}$  and  ${}^{1}A_{1} \rightarrow {}^{1}T_{1}$  of Fe<sup>II</sup>(LS), while Mössbauer spectra revealed a decrease in the Fe<sup>II</sup>(HS) peak. These results suggests that the observed photomagnetic phenomenon was induced by the photo-induced spin crossover from Fe<sup>II</sup>(LS) to Fe<sup>II</sup>(HS).

The photomagnetism of FeNb(4-pyoxm) complexes can be explained by a scheme described below. First, in the LT phase before light irradiation, a paramagnetic ion Nb<sup>IV</sup> and a diamagnetic ion Fe<sup>II</sup>(LS) are alternately bridged via a cyano group, and thus the spin source is physically remote, leading paramagnetism. Upon irradiation with a 473 nm light, Fe<sup>II</sup>-(LS), whose ground state is  ${}^{1}A_{1}$ , is excited to a singlet state  ${}^{1}T_{2}$ (or  ${}^{1}T_{1}$ ), a part of which subsequently transits to a quintet state,  ${}^{5}T_{2}$ , through a triplet state  ${}^{3}T_{2}$  or  ${}^{3}T_{1}$  (Figure 9).



**FIGURE 9.** Schematic illustration of light-induced excited spin-state trapping (LIESST) phenomenon.



**FIGURE 10.** Schematic illustration of the mechanism of ferromagnetic ordering due to photoinduced spin crossover in the FeNb(4-pyoxm) octacyano-complex.

In a photoinduced phase, which is metastable, the spins of photoinduced Fe<sup>II</sup>(HS) and those of neighboring Nb<sup>IV</sup> are antiferromagnetically aligned by the superexchange interaction via cyano groups, leading spontaneous magnetization (Figure 10). Note that the superexchange interaction constant ( $J_{ex}$ ) between Fe<sup>II</sup>(HS) and Nb<sup>IV</sup> is calculated by molecular field theory,<sup>39,49</sup> that is,  $|J_{ex}| = 3k_BT_{C}{Z_{Nb^{IV}Fe^{II}(HS)}Z_{Fe^{II}(HS)Nb^{IV}}S_{Fe^{II}(HS)}(S_{Fe^{II}(HS)}+1)S_{Nb^{IV}}(S_{Nb^{IV}}+1)}^{-1/2}$ , where the spin quantum numbers are  $S_{Nb^{IV}} = 1/2$  and  $S_{Fe^{II}(HS)} = 2$ , the numbers of the nearest neighbor are  $Z_{Nb^{IV}Fe^{II}(HS)} = 4$  and  $Z_{Fe^{II}(HS)Nb^{IV}} = 2$ ,  $k_B$  is the Boltzmann constant, and  $T_C = 20$  K. Consequently,  $J_{ex}$  was estimated to be -6.9 cm<sup>-1</sup>.

In the present study, we have synthesized a new 3D network metal complex  $Fe_2[Nb(CN)_8](4$ -pyridinealdoxime)\_8 $\cdot$  2H<sub>2</sub>O and succeeded in the observation of a photoinduced spin crossover ferromagnetism for the first time. This photomagnetic effect is induced by the LIESST and the

superexchange interaction between Fe<sup>II</sup>(HS) and Nb<sup>IV</sup>. In the previous work, we reported Fe<sub>2</sub>[Nb(CN)<sub>8</sub>]·(3-pyCH<sub>2</sub>OH)<sub>8</sub>·  $4.6H_2O$ ,<sup>25</sup> in which spin-crossover and ferrimagnetism coexist, but this compound did not show the light-induced spin-crossover ferromagnetism. This is because the crystal structure is very highly symmetric (cubic:  $Ia3\overline{d}$ ) and results in a very stable alternating arrangement of  $\cdots$  –Fe<sup>II</sup>(HS)– Fe<sup>II</sup>(LS)– $\cdots$ , and consequently, photoinduced Fe<sup>II</sup>(HS) domains cannot be formed. Hence, a somewhat lowsymmetrical crystal structure is effective to produce photoinduced Fe<sup>II</sup>(HS) domains, which is necessary to realize LIESST-induced bulk magnetization.

#### 6. Conclusions

In addition to the materials described herein, we are conducting research on the photoinduced magnetic pole inversion in a hexacyanometal complex system, (Fe<sup>III</sup><sub>0.4</sub>- $Mn^{II}_{0.6})_{1.5}[Cr^{III}(CN)_6] \cdot 7.5H_2O^{38,39}$  and photomagnetism in octacyanometal complex systems, Cs<sup>1</sup><sub>2</sub>Cu<sup>II</sup><sub>7</sub>[Mo<sup>IV</sup>(CN)<sub>8</sub>]<sub>4</sub>.  $6H_2O_{1}^{54}$  [Cu<sup>II</sup>(1,4,8,11-tetraazacyclodecane)]<sub>2</sub>[Mo<sup>IV</sup>(CN)<sub>8</sub>].  $10H_2O_{,56}^{56}$  and  $Cs^{I}Co^{II}(3-cyanopyridine)_2[W^{V}(CN)_8] \cdot H_2O_{,50}^{50}$ In hexacyanometal complexes, the photomagnetic phenomena of a K<sub>0.2</sub>Co<sub>1.4</sub>[Fe(CN)<sub>6</sub>]·6.9H<sub>2</sub>O system and the photoinduced change of magnetic susceptibility in a  $K_{1.54}V^{II}_{0.77}V^{III}_{0.08}[Cr^{III}(CN)_6](SO_4)_{0.16} \cdot 3.1H_2O$  system are known.<sup>42,82</sup> In systems where photochromic molecules are introduced into magnetic complexes such as oxalate metal complexes and single-molecule magnets, changes in magnetic coercive field and alternating current (ac) magnetic susceptibility induced by photoisomerization are observed.83-85

Besides these molecular magnetic materials, photomagnetic effects are also reported in dilute magnetic semiconductors and perovskite-type manganese oxides.<sup>86,87</sup> Furthermore, theoretical studies on photoinduced phase transition are actively being conducted.<sup>88–93</sup> Metal complexes are superior as a model material to explore novel functional materials. By taking advantage of the knowledge that has been obtained through our previous studies, we have recently discovered a novel metal oxide,  $\lambda$ -trititanium pentoxide ( $\lambda$ -Ti<sub>3</sub>O<sub>5</sub>), which exhibits photoreversible metal–semiconductor phase transition at room temperature,<sup>94</sup> suggesting the possibility of developing room-temperature photomagnetism by concurrently studying magnetism.

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#### FOOTNOTES

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